

ASSOCIATIVE PHOTOPRODUCTION OF ROPER RESONANCE WITH VECTOR MESONS

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ABSTRACT

We examine the associative photoproduction of Roper resonance with vector mesons, $\gamma+N \rightarrow N^*(1440)+V$, in the near threshold region by using the analogy with the processes of “elastic” vector meson photoproduction, $\gamma+N \rightarrow N+V$. Different production mechanisms are suggested. In addition to standart σ - and π -exchanges in the t-channel, the baryonic exchanges in the s- and u-channel are taken into account as well. We calculate the differential cross-section and beam asymmetry on proton and neutron targets for the reaction $\gamma+N \rightarrow N^*(1440)+\omega$. All calculations are performed at $E_\gamma = 2.5$ GeV.

The recent study [1] on the associative photoproduction of Roper resonance with vector mesons, $\gamma+N \rightarrow N^*(1440)+V$, in the near threshold region indicates the importance of the pseudoscalar (π) exchange for t-channel in associative ω -meson photoproduction but it does not give a clear information about the relative role of different possible threshold mechanisms on the production of $N^*(1440)$. Therefore, in additon to standart σ - and π -exchanges in the t-channel, we consider the baryonic resonances in the s- and u-channel .

For t-channel, the pseudoscalar (π) and scalar exchanges (σ), Fig.1(a), are considered The pseudoscalar exchange amplitude can be obtained from the Lagrangian,

$$\mathbf{L}_\pi = g_{\omega\pi\gamma} \varepsilon^{\mu\nu\alpha\beta} \partial_\mu V_\nu \partial_\alpha A_\beta \pi - i g_{\pi NN^*} \bar{N}^* \gamma_5 \pi N \quad (1)$$

where A_μ is the photon field. Following Ref. [1,2], we use the coupling constants as $g_{\omega\pi\gamma}^2 = 3.315$ and $g_{\pi NN^*}^2/4\pi = 3.4$. The form factors that we have used are as follows:

$$F_{\pi NN^*} = \frac{\Lambda_\pi^2 - M_\pi^2}{\Lambda_\pi^2 - t}, \quad F_{\omega\pi\gamma} = \frac{\Lambda_{\omega\pi\gamma}^2 - M_\pi^2}{\Lambda_{\omega\pi\gamma}^2 - t}, \quad (2)$$

where $\Lambda_\pi = 0.7$ and $\Lambda_{\omega\pi\gamma} = 0.77$ in GeV unit. [3]

The scalar (σ) exchange amplitude can be obtained from the Lagrangian,

$$\mathbf{L}_\sigma = \frac{e g_{\omega\sigma\gamma}}{M_\omega} (\partial^\mu V^\nu \partial_\mu A_\nu - \partial^\mu V^\nu \partial_\nu A_\mu) \sigma + g_{\sigma NN^*} \bar{N}^* N \sigma \quad (3)$$

The form factors for this exchange are given as,

$$F_{\sigma NN^*} = \frac{\Lambda_\sigma^2 - M_\sigma^2}{\Lambda_\sigma^2 - t}, \quad F_{\omega\sigma\gamma} = \frac{\Lambda_{\omega\sigma\gamma}^2 - M_\sigma^2}{\Lambda_{\omega\sigma\gamma}^2 - t} \quad (4)$$

where $\Lambda_\sigma = 1.0$ GeV, $\Lambda_{\omega\sigma\gamma} = 0.9$ GeV, $M_\sigma = 0.5$ GeV and the coupling constant, $g_{\sigma NN}^2/4\pi=0.34$ [1,2].

Finally, the baryonic resonances in the s- and u-channel, Figure 1 (b,c), are considered. The Lagrangian for the (s+u) nucleonic resonances can be given by

$$\mathbf{L}_S = \bar{N}^* \left[g_{\omega NN^*}^V \gamma_\mu V^\mu - \frac{g_{\omega NN^*}^T}{M_N + M^*} \sigma_{\mu\nu} \partial^\nu V^\mu \right] N - e \bar{N}^* \left[Q_N \gamma_\mu A^\mu - \frac{\kappa_N}{2M_N} \sigma_{\mu\nu} \partial^\nu A^\mu \right] N \quad (5)$$

$$\mathbf{L}_U = \bar{N}^* \left[g_{\omega NN^*}^V \gamma_\mu V^\mu - \frac{g_{\omega NN^*}^T}{M_N + M^*} \sigma_{\mu\nu} \partial^\nu V^\mu \right] N - e \bar{N}^* \left[Q_N \gamma_\mu A^\mu - \frac{\kappa_{N^*}}{2M^*} \sigma_{\mu\nu} \partial^\nu A^\mu \right] N$$

where κ_N is the anomalous magnetic moment of the nucleon ($\kappa_p = 1.79$, $\kappa_n = -1.91$), $g_{\omega NN^*}^T$ and $g_{\omega NN^*}^V$ are the tensor and vector coupling constants, respectively. κ_{N^*} is the anomalous magnetic moment of the $N^*(1440)$, and, for the simplicity, in our calculations we take

$$\kappa_{N^*} = \kappa_N, \quad g_{\omega NN^*}^V = 0$$

The differential cross-section and beam asymmetry for the Roper resonance, $\gamma + N \rightarrow N^*(1440) + \omega$, are calculated at $E_\gamma = 2.5$ GeV. The values of the coupling constants, $g_{\omega NN^*}^T$ and $g_{\omega\sigma\gamma}$, are 3.41 and 0.13. The value of the latter coupling constant is taken from reference [4,5]. The results for our model calculations are shown in Figure 2. From the curves, it can be concluded that the σ -exchange contribution is so small that the behaviour of the graphs for π and $(\pi+\sigma)$ terms does not change in both differential cross-section and beam asymmetry, and, the dominance of the s+u nucleonic term contribution is obvious as the value of $|t|$ (momentum transfer) increases.

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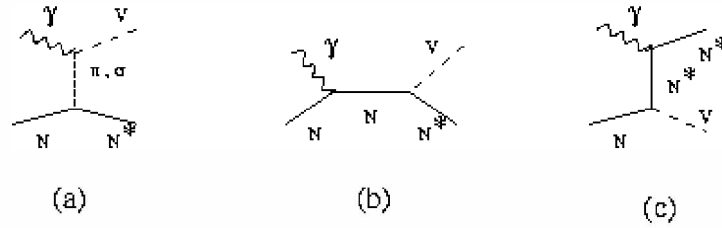


Figure 1: Three mechanisms for vector meson ($V = \omega$) photoproduction involving the excitation of the $N^*(1440)$ resonance: (a) one-boson exchange, (b) and (c) s- and u-channel intermediate nucleon diagrams.

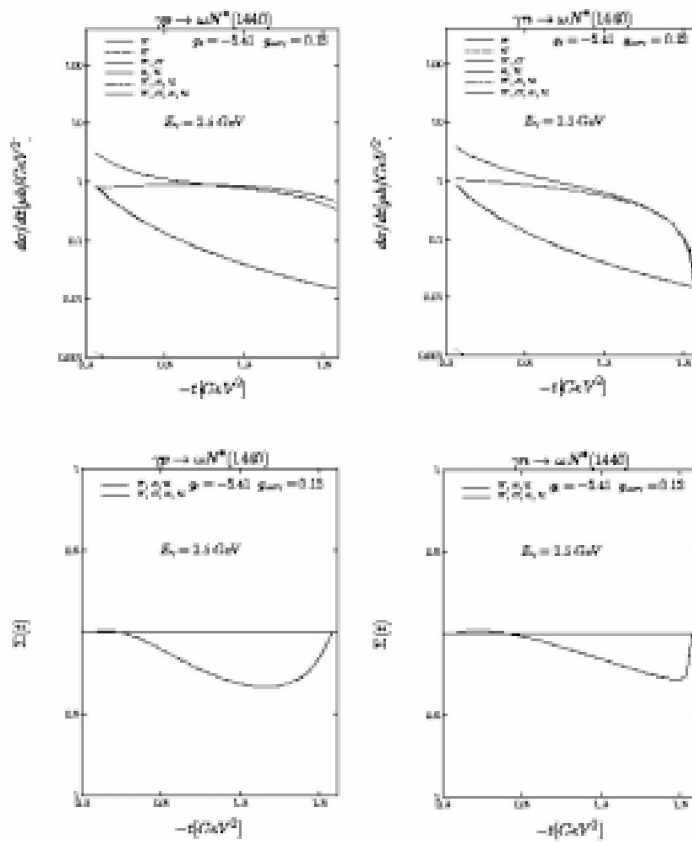


Figure 2 : Differential cross-section and beam asymmetry for inelastic ω meson photoproduction on proton and neutron targets at $E_\gamma = 2.5$ GeV